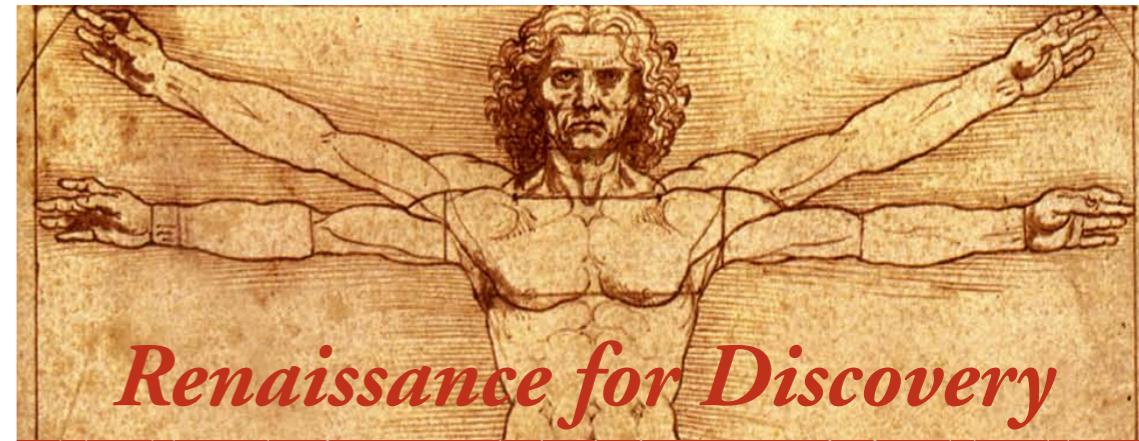


PHENIX HBT Measurements

Jason Newby (LLNL) 
for the PHENIX Collaboration



RHIC AGS Users' Meeting 2006

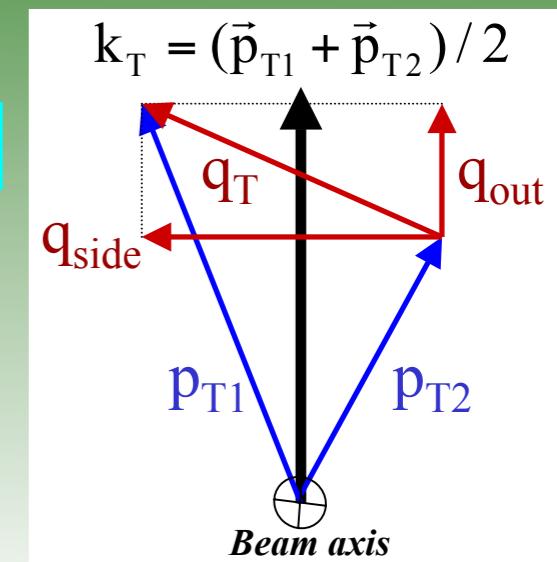
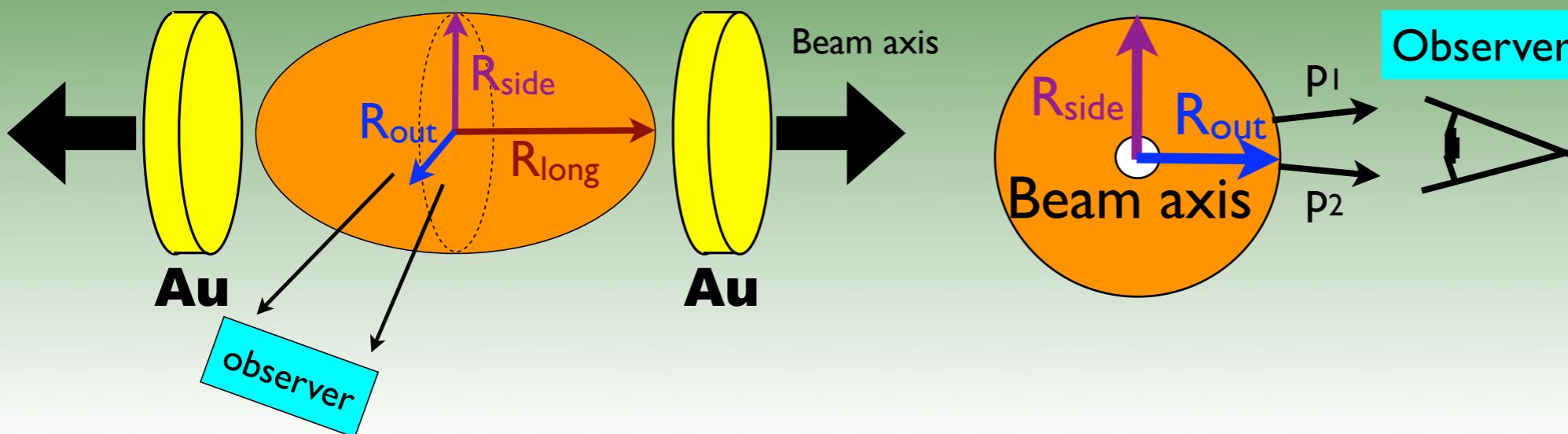


What does HBT offer?

- Femtoscopy utilizes Quantum Statistical (BEC) or Final State Interactions (Coulomb, Strong) that dominate multi-particle correlations at small separations in phase space to measure space-time structure.
- Multi-particle Femtoscopy is the most direct measurement we have of the space-time evolution of Heavy Ion Collisions at the freeze-out surface.
- Assuming Gaussian profiles, 1D and 3D radii are extracted from two-particle correlations in momentum difference distributions.
- Emerging Imaging Techniques make no assumptions of the profile offering sensitivity to finer structure.

HBT in Heavy-ions

Bertsch-Pratt parameterization



Bowler-Sinyukov

Phys. Lett. B270,69 (1991)

Phys. Lett. B432,249 (1998)

$$C_2(\vec{q}) = (1 - \lambda) - \lambda \exp \left[-R_{side}^2 q_{side}^2 - R_{out}^2 q_{out}^2 - R_{long}^2 q_{long}^2 - R_{so}^2 q_{side} q_{out} \right] K_{Coul}(q_{inv})$$

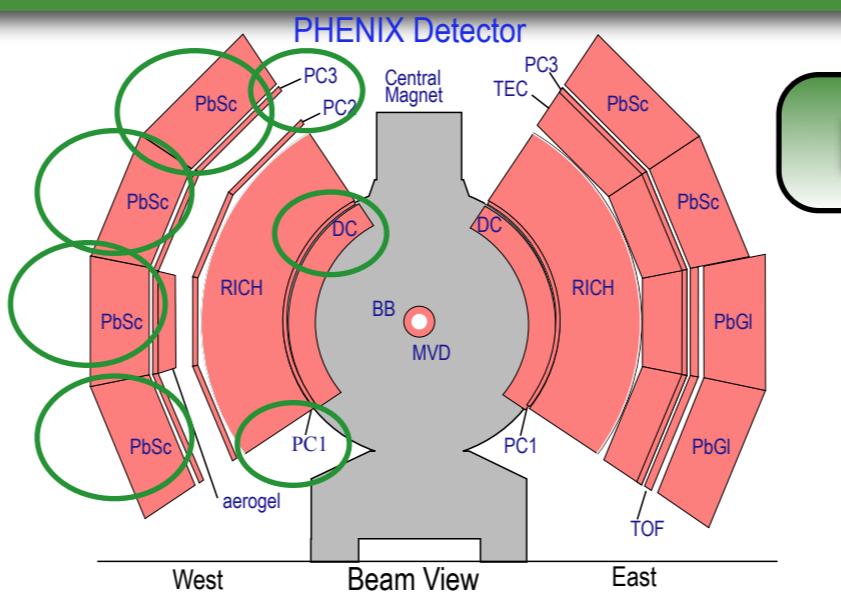
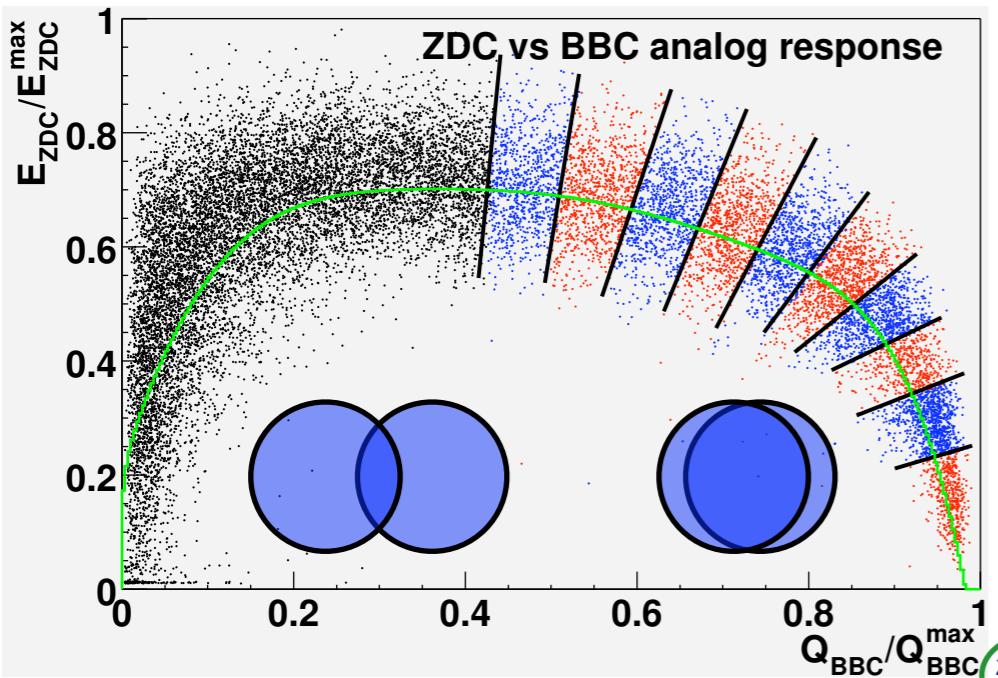
Fraction of Chaotic Emission

Coulomb Interaction

Halo of long lived resonances

PHENIX Detector

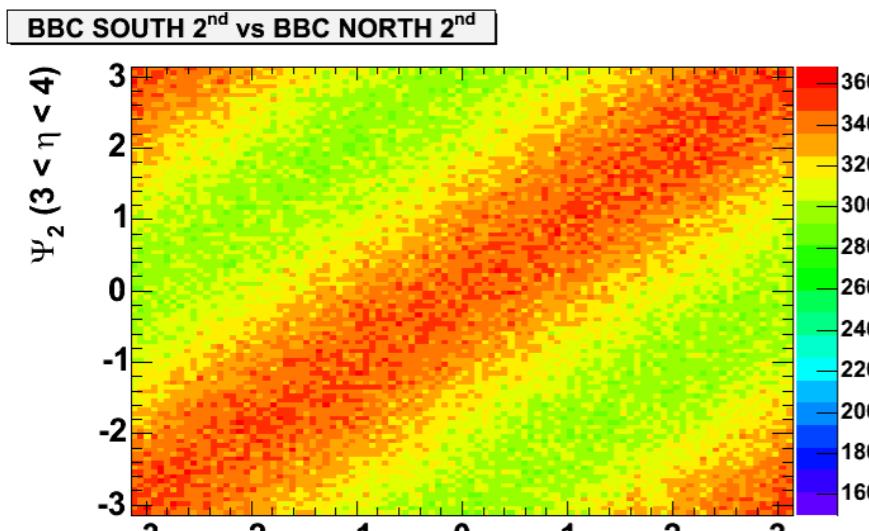
BBC/ZDC Centrality



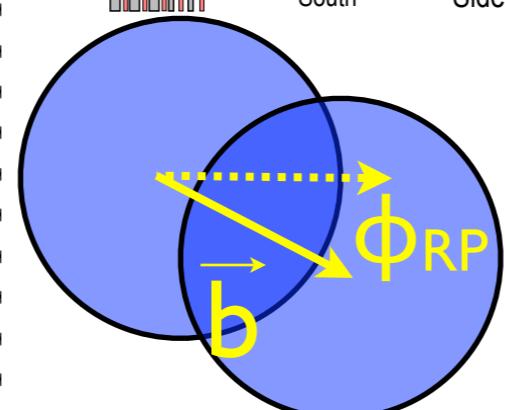
DC/PC Track-Mom Reco

$$\delta p/p \sim 1\%$$

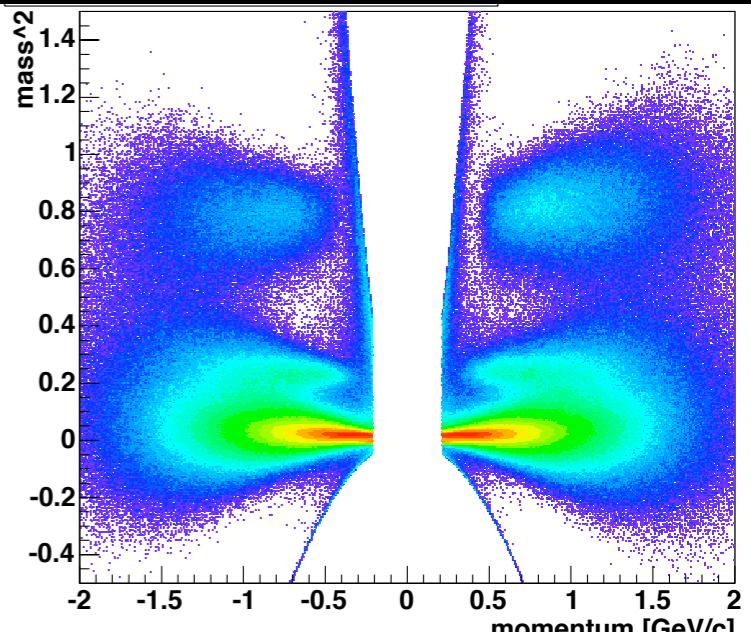
BBC Reaction Plane



$$\tan 2\Phi_{RP} = \frac{\sum n_{ch} \sin 2\phi_{PMT}}{\sum n_{ch} \cos 2\phi_{PMT}}$$



EMC Particle Identification



PHENIX HBT Overview

- PHENIX HBT analyses of $K^{+/-}, \pi^{+/-}, p$:
- Gaussian 1D/3D Radii from C_2
 - ▶ Full/Partial Coulomb Correction
 - ▶ Bowler-Sinyukov Core-Halo
 - ▶ Full Koonin-Pratt
- Levy Fits to 2/3 Particle Correlations
- 1D Source Imaging

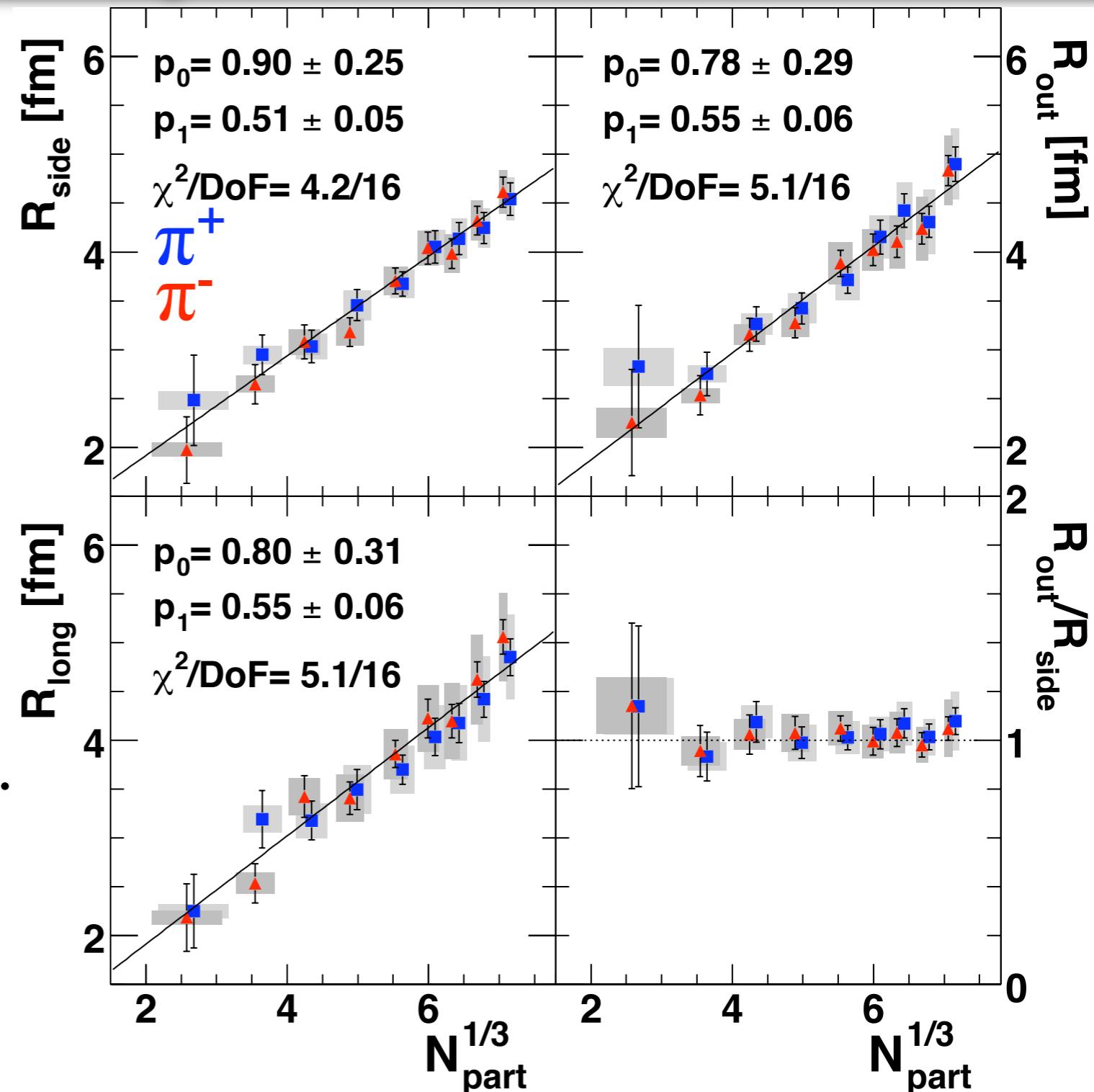
	$\pi^{+/-}$	$K^{+/-}$	p
Au+Au 130	3D (k_T)		
Run 2 Au+Au 200	3D (k_T, N_p)	ID	ID
Run 4 Au+Au 200	AzHBT ID Image	3D ID Image	



System Size Dependence

200 GeV Au+Au
PRL93(2004)152302

- Larger initial systems “appear” larger
- $N_p^{1/3}$ scaling of the pion source
- HBT Radii are dominated by geometric length scales.
- Expected long emission duration would increase $R_{\text{out}}/R_{\text{side}}$: not observed.



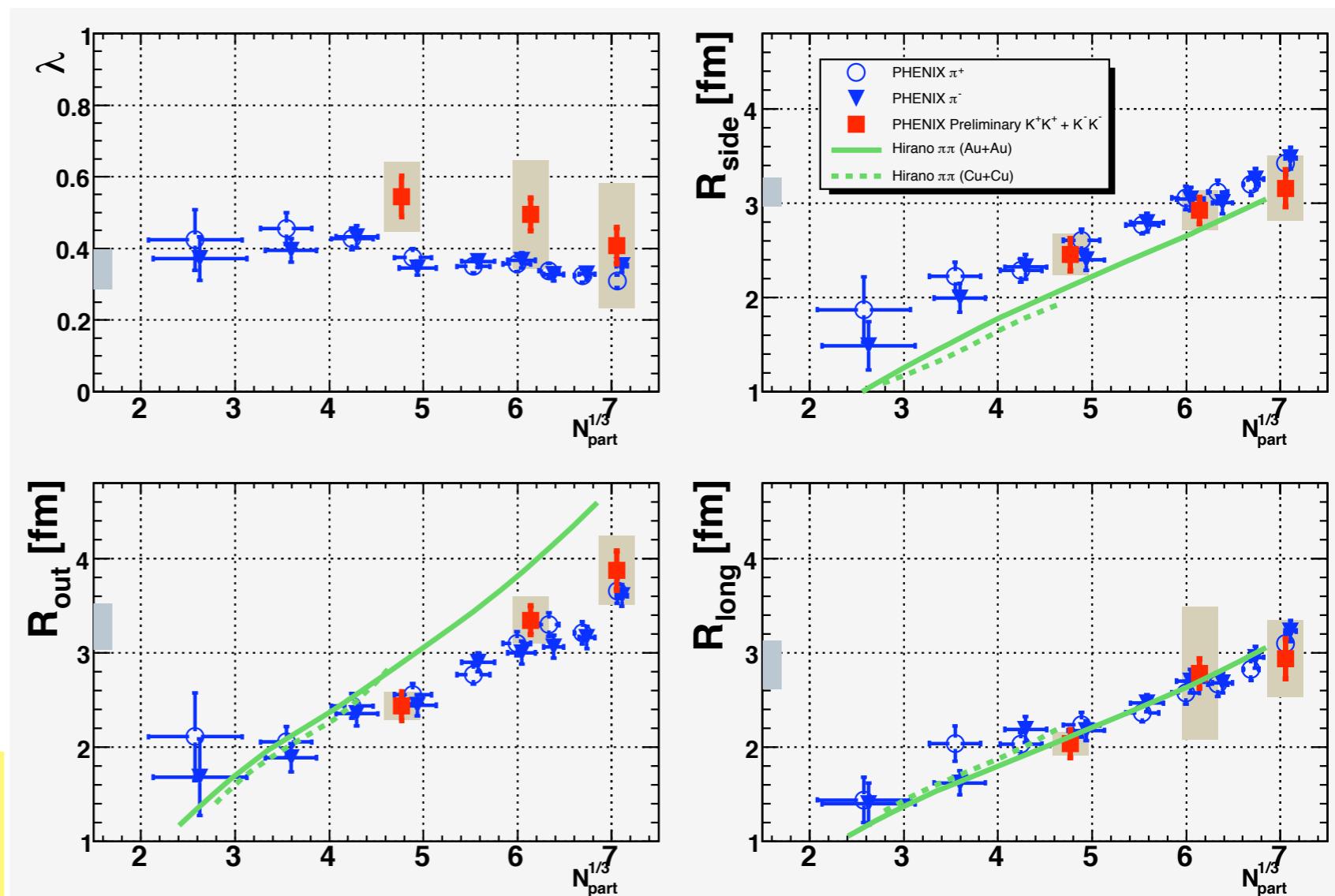
Kaon Participant Scaling

Au+Au 200 GeV

Pion $\langle m_T \rangle = 0.47$
Kaon $\langle m_T \rangle = 0.89$

R_{out} , R_{side} scaled to
match kaon $\langle m_T \rangle$

Hydro agreement with
data is quite good.



A. Enokizono, DNP Maui 2005
(T.Hirano, Y.Nara, Partial Chemical Freeze-out,
NPA 743('04)305)



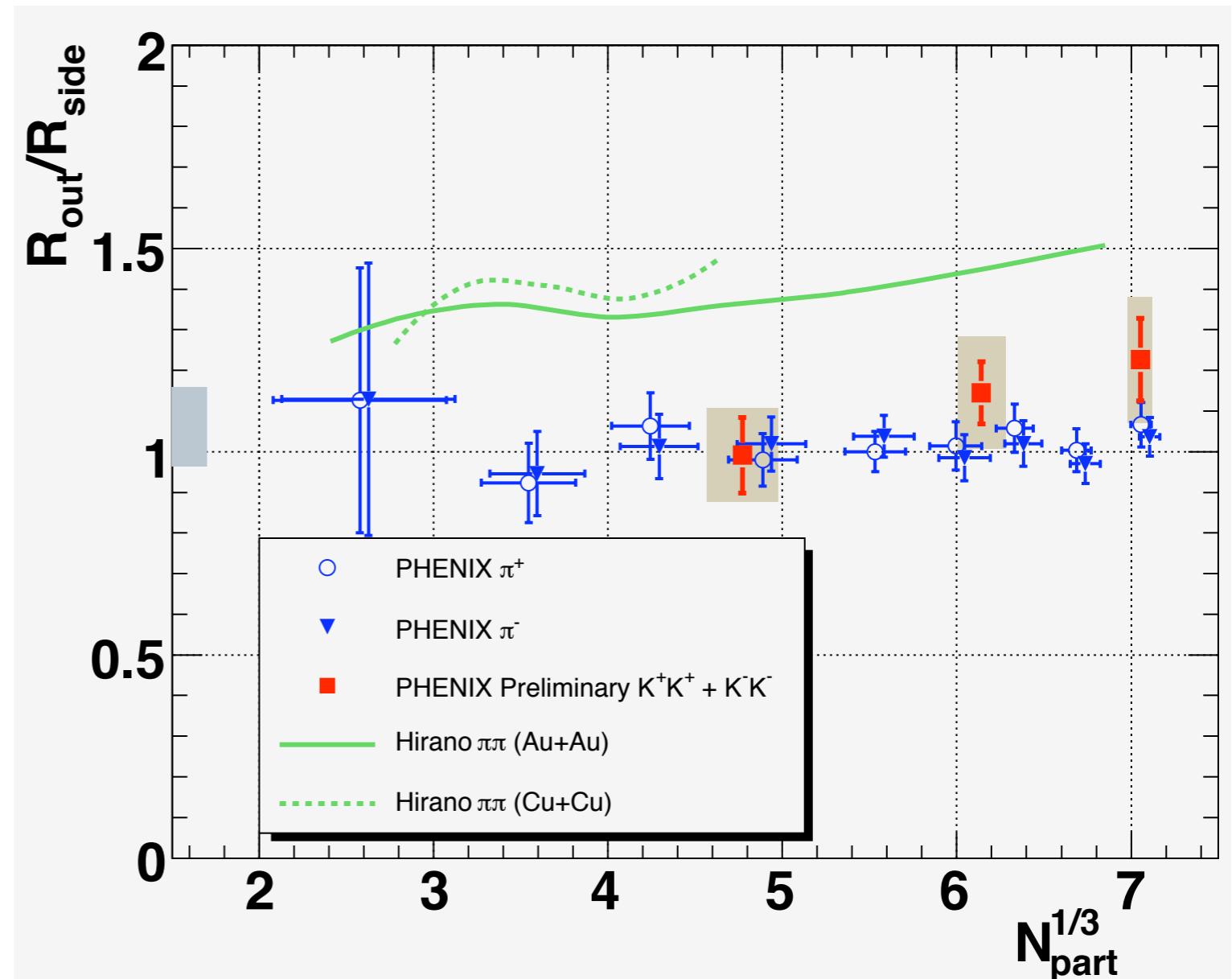
Kaon $R_{\text{out}}/R_{\text{side}}$

Au+Au 200 GeV

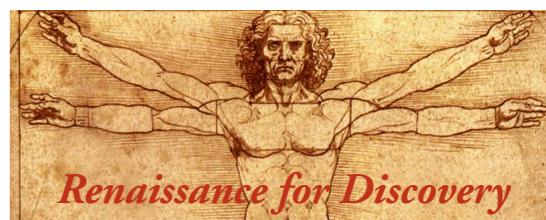
Pion $\langle m_T \rangle = 0.47$
Kaon $\langle m_T \rangle = 0.89$

$R_{\text{out}}, R_{\text{side}}$ scaled to
match kaon $\langle m_T \rangle$

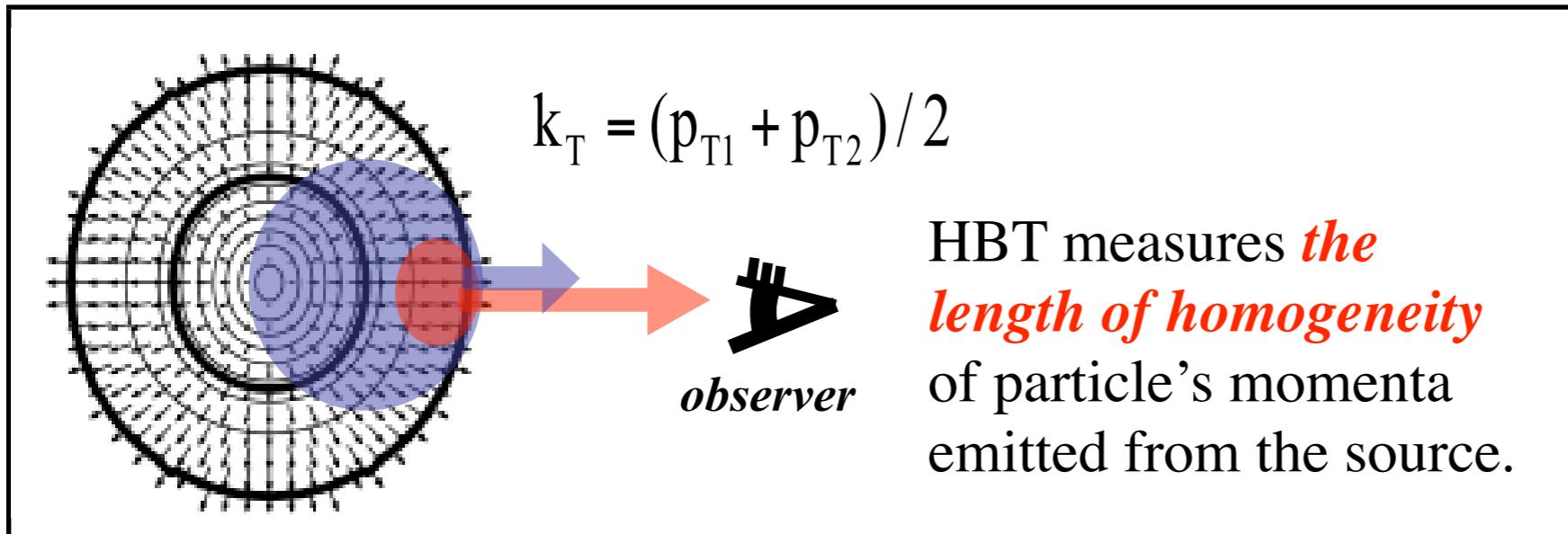
“Hint” of $R_{\text{out}}/R_{\text{side}} > 1$
for central collisions.



A. Enokizono, DNP Maui 2005
(T.Hirano, Y.Nara, Partial Chemical Freeze-out,
NPA743('04)305)



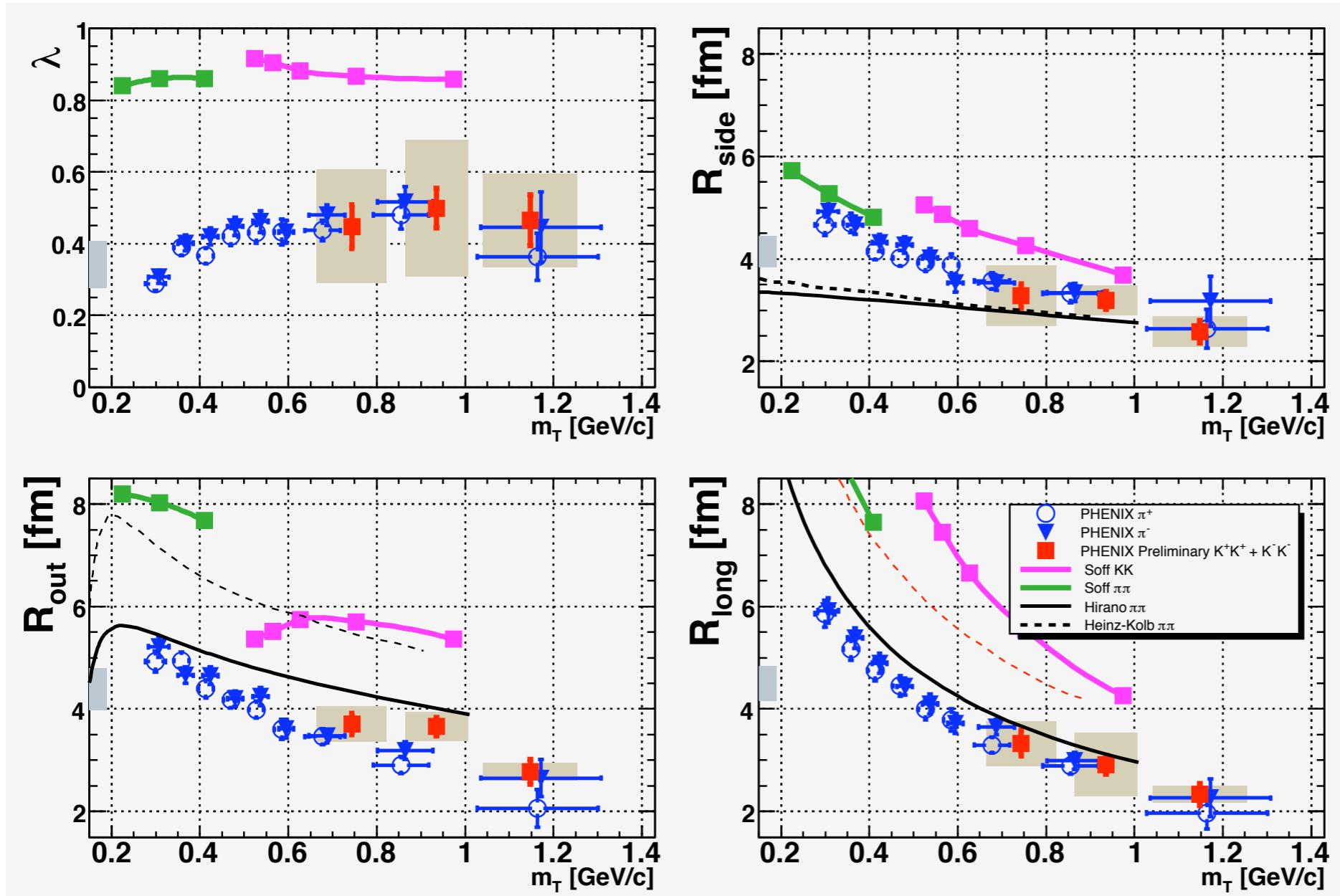
x-p Correlations



- HBT actually measures the space-time structure of “regions of homogeneity”
- Relativistic Expansion reduces the length of homogeneity for high momentum pairs.

$$R_{long} \propto \frac{1}{\sqrt{m_T}}$$

Kaon m_T Dependence



A. Enokizono, DNP Maui 2005

Hydro+UrQMD $T_c=160\text{MeV}$

(S. Soff, hep-ph/0202240)

3D Hydro

Hirano, PRC66(2002)054905

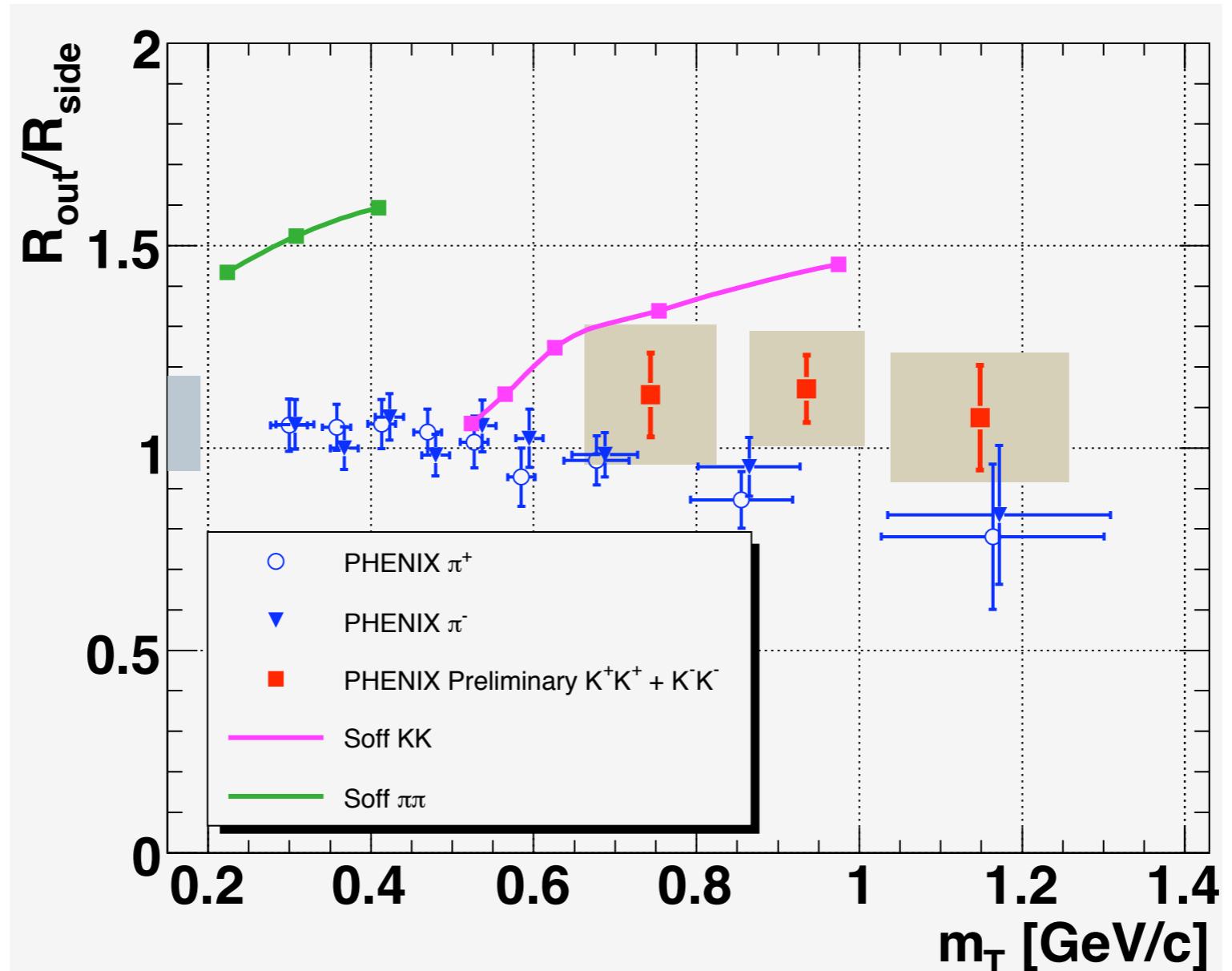
2D+1 Hydro

Heinz-Kolb, hep-ph/0204061

Small rescattering in hadronic phase.



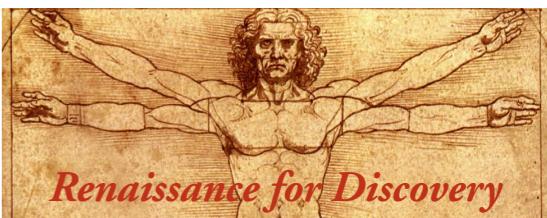
$R_{\text{out}}/R_{\text{side}}$ m_T Dependence



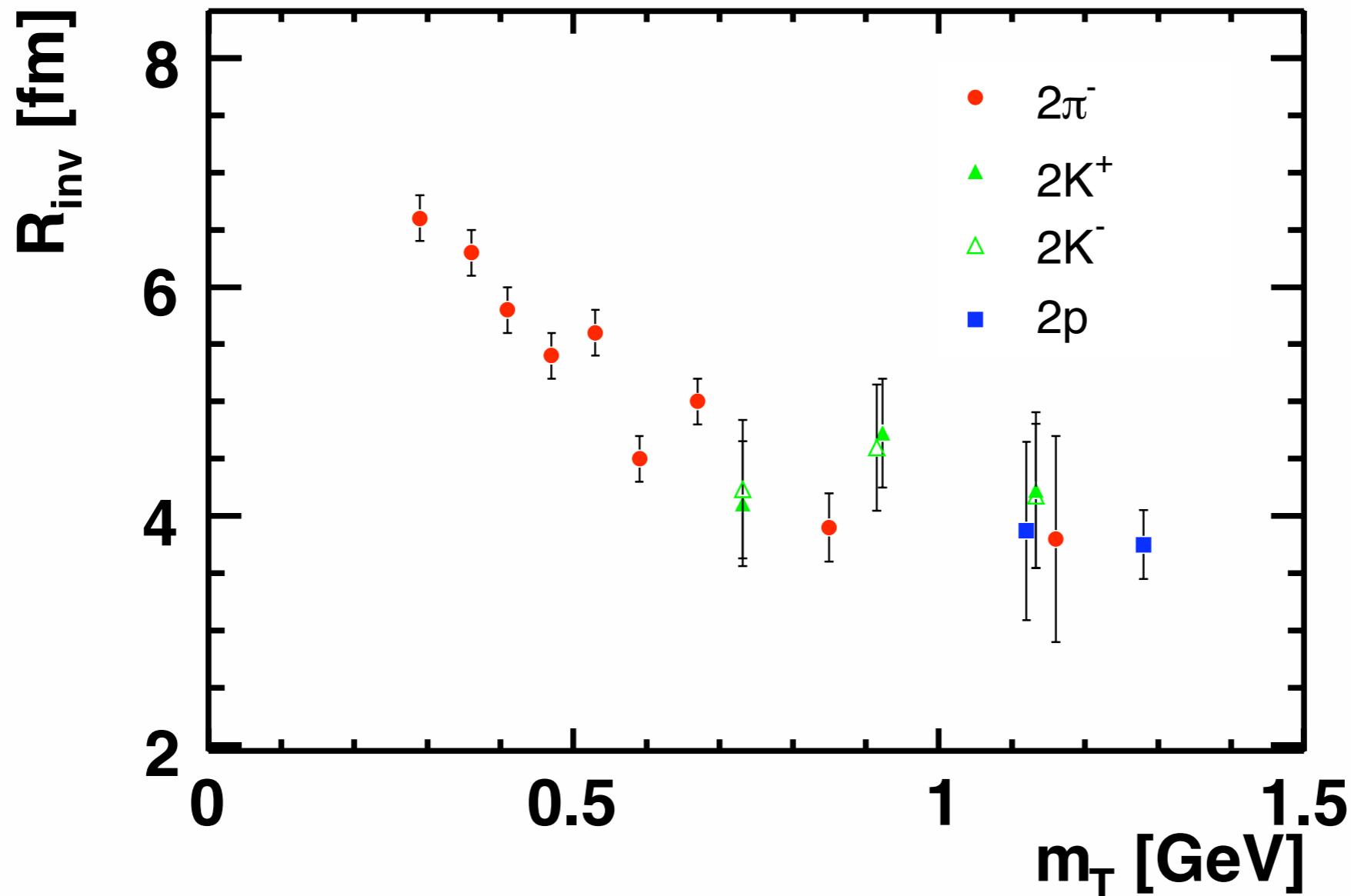
A. Enokizono, DNP Maui 2005

Hydro+UrQMD (S. Soff, hep-ph/0202240, Tc=160MeV)

Kaons give no indication of long emission duration



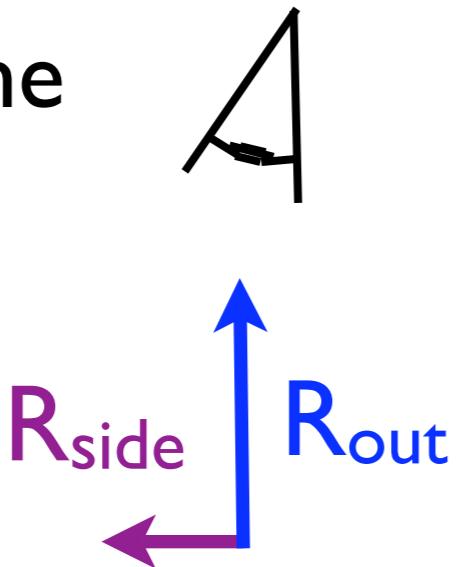
ID Species Dependence



Protons follow the pion and kaon trend.
No species dependence observed.

Azimuthal HBT

Out-of-Plane
Observer
 $\phi = 90^\circ$



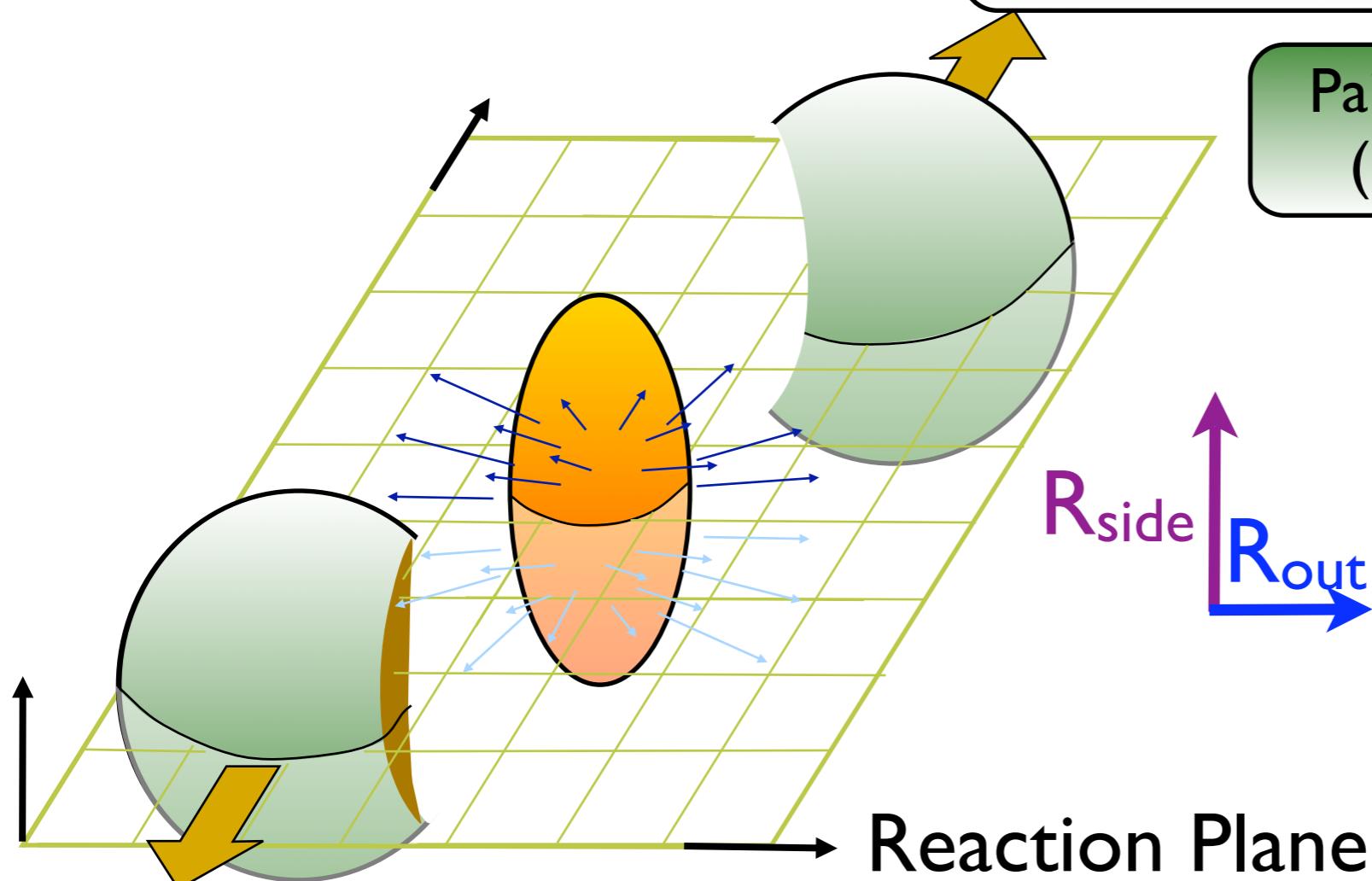
Heinz et al. PRC66,044903(2002)

$$R_\mu^2(\Phi) = R_{\mu,0}^2 + R_{\mu,n}^2 \cos(n\Phi) \quad \{\mu=s,o,l\}$$

$$R_{s,2}^2 > 0 \quad R_{o,2}^2 < 0$$

$$R_{so}^2(\Phi) = R_{so,0}^2 + R_{so,n}^2 \sin(n\Phi)$$

$$R_{os,2}=0 \quad R_{os,2}>0$$



Pair Ordering
($p_{z,2} > p_{z,1}$)

Pairs Boosted
LCMS ($p_{z,1} = -p_{z,2}$)

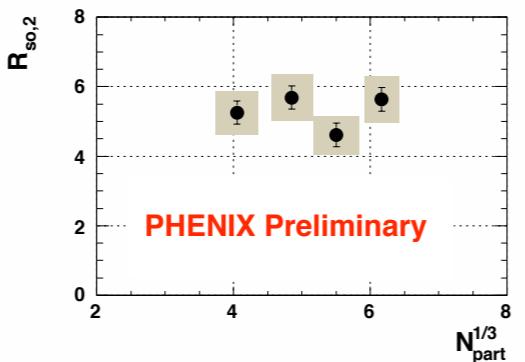
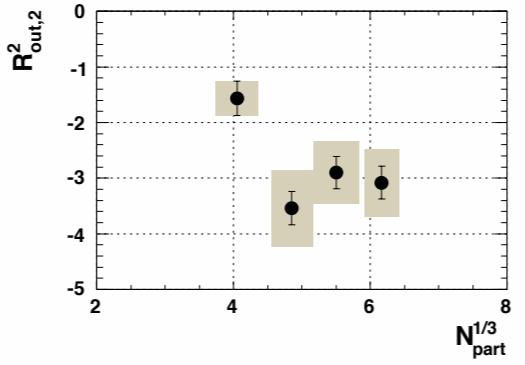
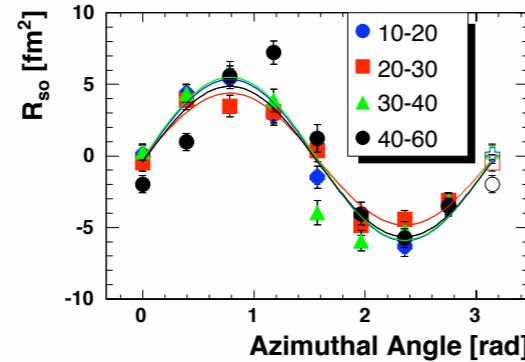
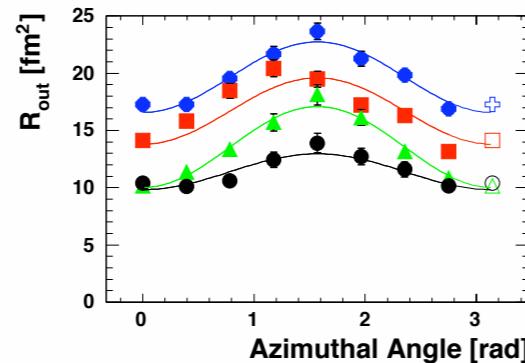
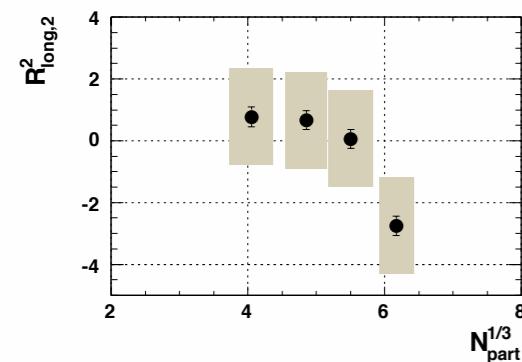
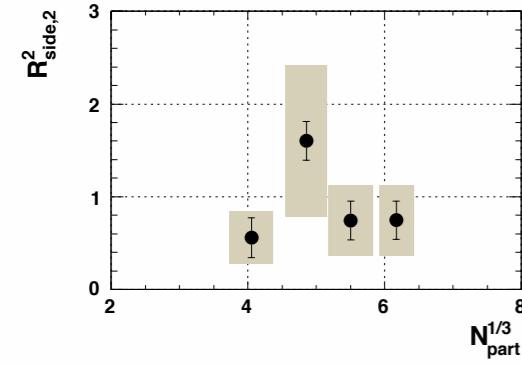
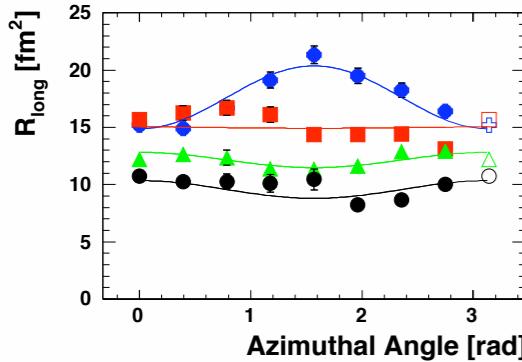
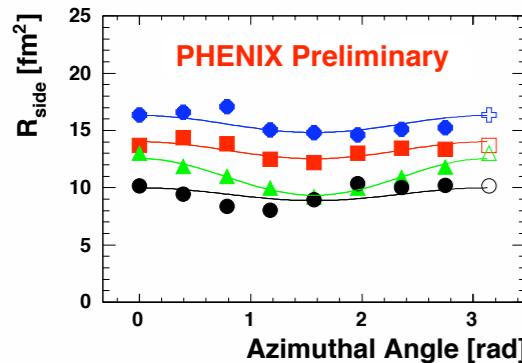
$$\varepsilon = \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle} \approx \frac{2R_{s,2}^2}{R_{s,0}^2}$$

In-Plane
Observer
 $\phi = 0^\circ$

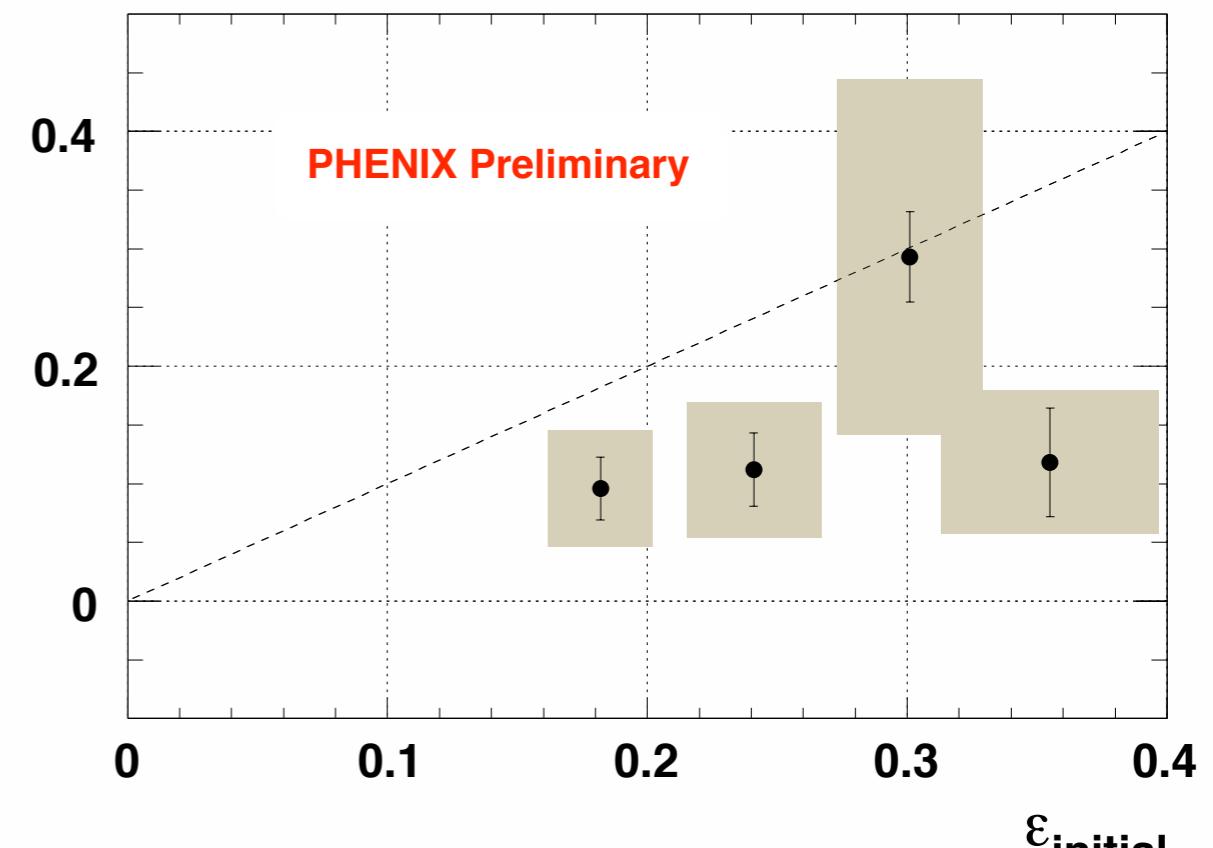


Renaissance for Discovery

Eccentricity Final State vs Initial

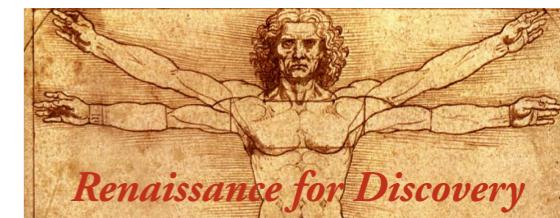


$$\epsilon \approx 2R_{s,2}/R_{s,0}$$



Final state eccentricity $\sim \frac{1}{3}$ to $\frac{1}{2}$ of initial eccentricity.

Reduction of systematic errors in progress.



Renaissance for Discovery

Testing the Core-Halo

- λ from $\pi\pi$ C_2 and C_3
- Constrain coherent and extended source contributions

hep-ph/0001233, Nucl. Phys. B329 (1990) 357

$$f_c(p) = N_c(p)/N_1(p)$$

$$p_c(p) = N_c^p(p)/N_c(p)$$

$N_1(p)$... Invariant mom. distr.

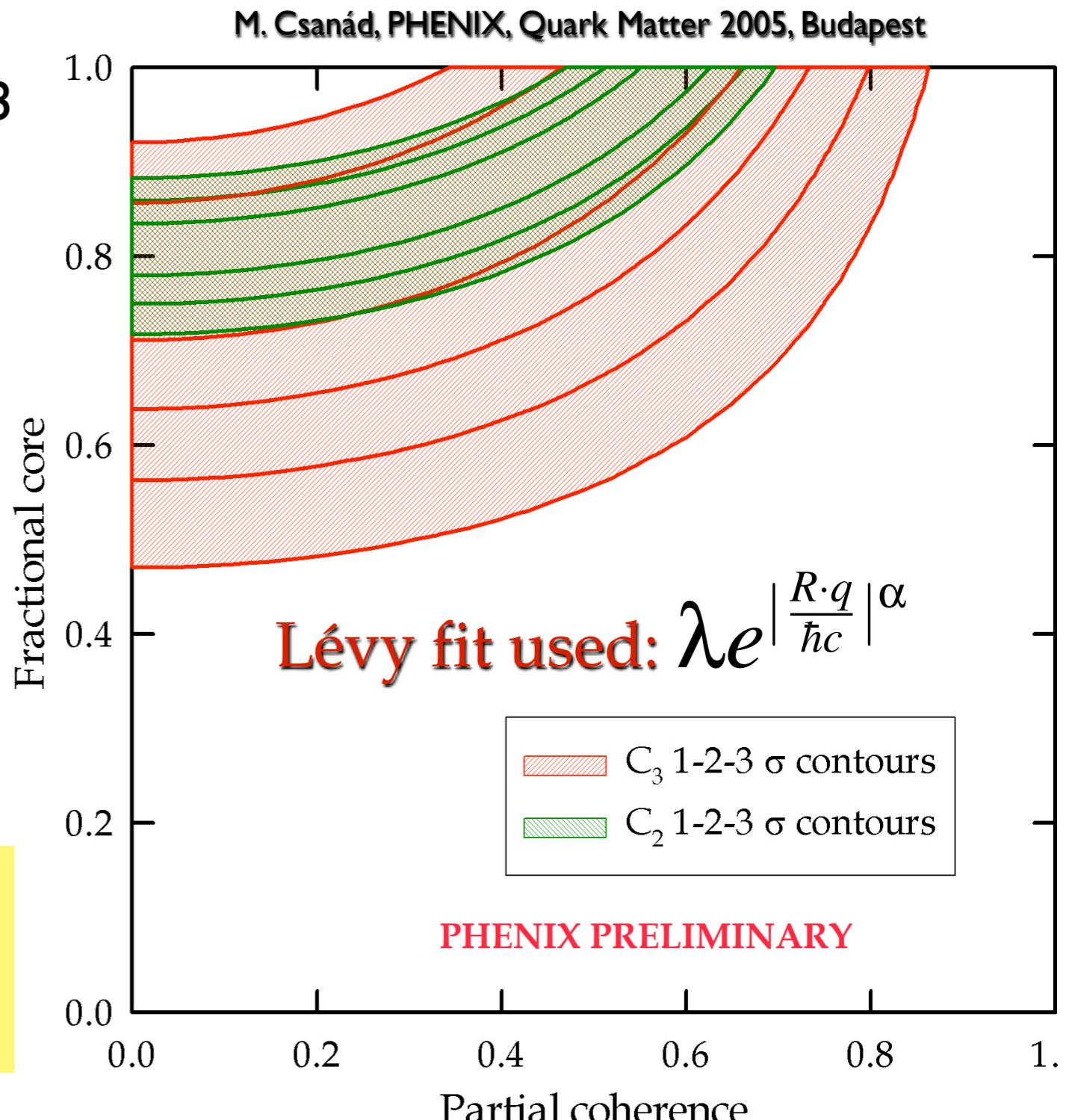
$N_c(p)$... Core fraction

$N_c^p(p)$... Part. coh. fraction

λ alone - unable to unambiguously distinguish resonance contribution vs extended source

$$C_2(p_1 \simeq p_2) = 1 + f_c^2[(1 - p_c)^2 + 2p_c(1 - p_c)]$$

$$C_3(p_1 \simeq p_2 \simeq p_3) = 1 + 3f_c^2[(1 - p_c)^2 + 2p_c(1 - p_c)] + 2f_c^3[(1 - p_c)^3 + 3p_c(1 - p_c)^2]$$



Source Imaging

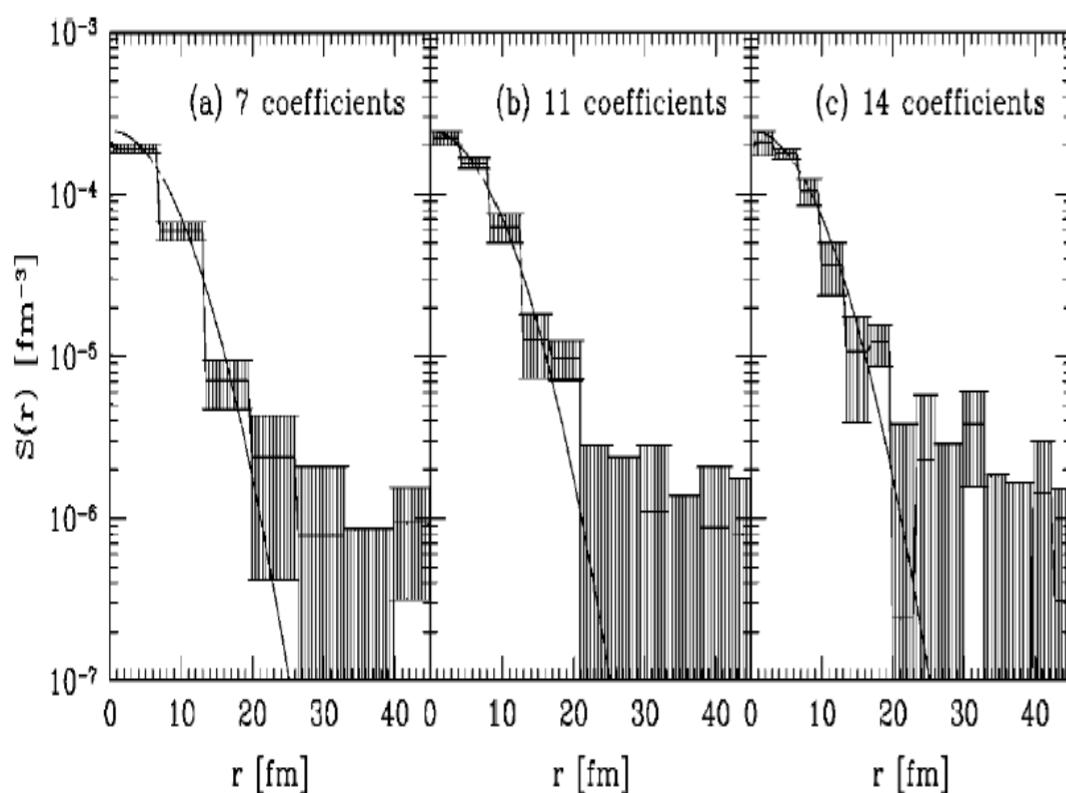
$$R_{\vec{P}}^{\text{obs}}(\vec{q}) \equiv C_{\vec{P}}^{\text{obs}}(\vec{q}) - 1 = \int d\vec{r} K(\vec{q}, \vec{r}) S_{\vec{P}}(\vec{r})$$

$$K(\vec{q}, \vec{r}) = |\Phi_{\vec{q}}(\vec{r})|^2 - 1$$

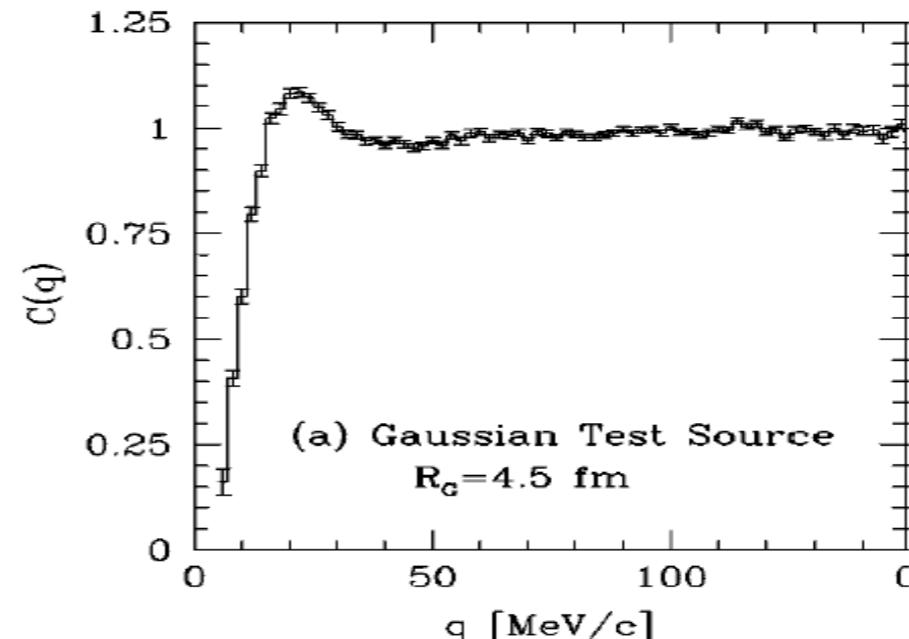
is kernel which can be calculated from BEC and known final state interactions of pairs.

$$S_{\vec{P}}(\vec{r})$$

is source function which represents the emission probability of pairs at r in the pair CM frame.



D.A. Brown and P. Danielewicz
Phys. Rev. C. 64, 014902 (2001)



$S(r)$ is expanded in a function basis.

$$S(r) = \sum_{i=1}^{N_M} S_i B_i(r) \quad \Delta S(r) = \sqrt{\sum_{i,j=1}^{N_M} \Delta^2 S_{ij} B_i(r) B_j(r)}$$

Imaging by minimization of χ^2

$$\chi^2 = \sum_i \left(R_i^{\text{obs}} - \sum_j K_{ij} S_j \right)^2 / \sigma^2 R_{ii}^{\text{obs}}$$

Most probable source function is

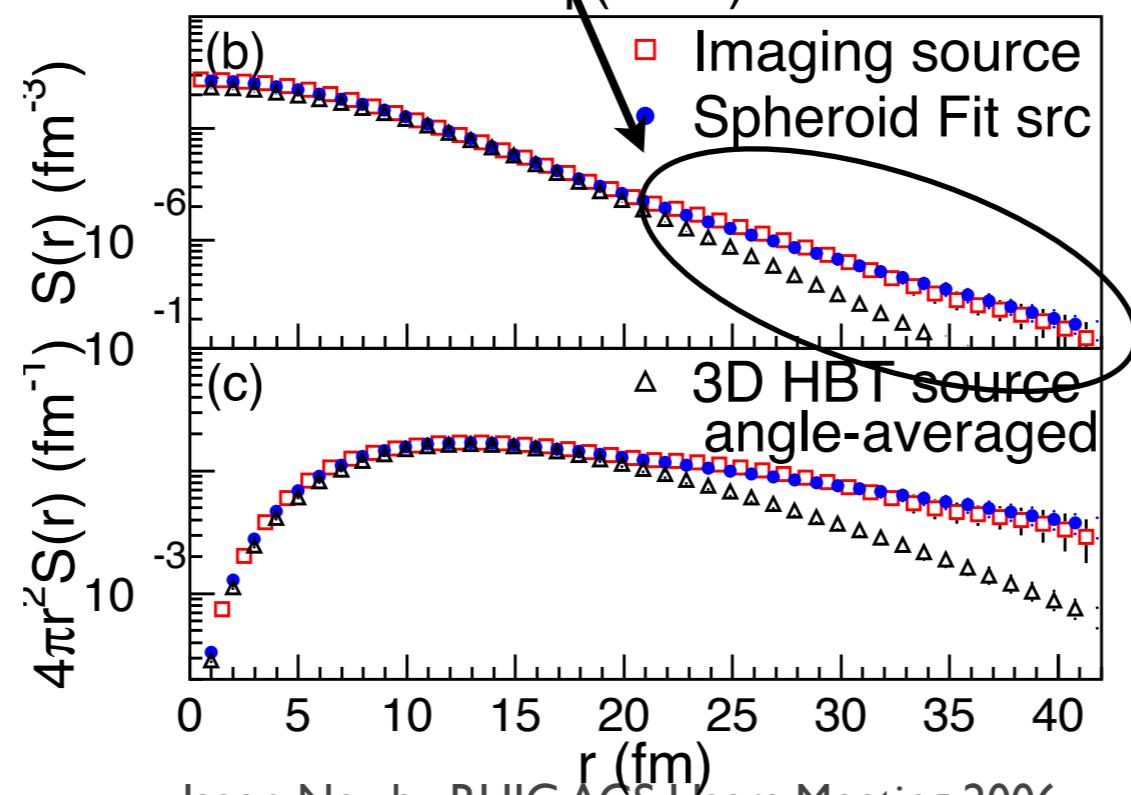
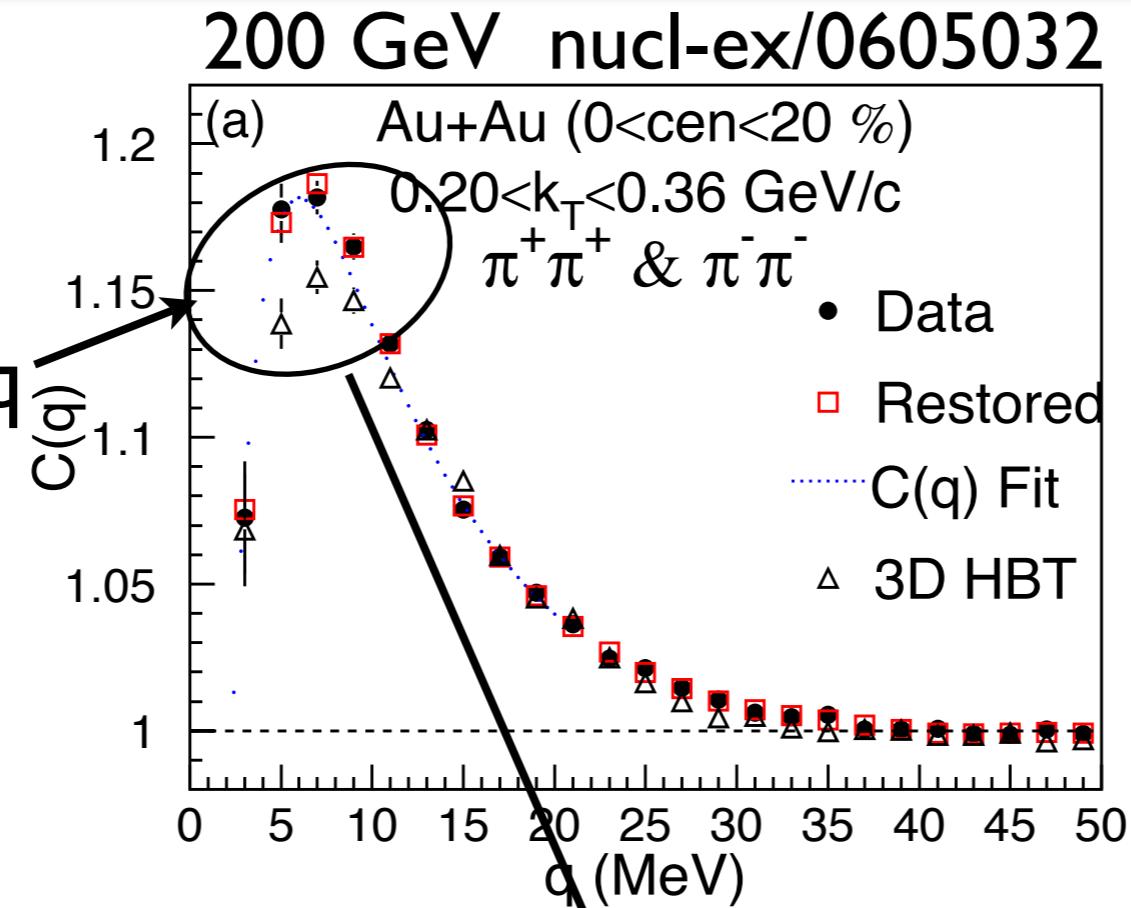
$$S = \hat{A}^2 S \times K^T \times (\hat{A}^2 R^{\text{obs}})^{-1} \times R^{\text{obs}}$$

$$\hat{A}^2 S = \left(K^T \times (\hat{A}^2 R^{\text{obs}})^{-1} \times K \right)^{-1}$$



Measured 1D Source Image

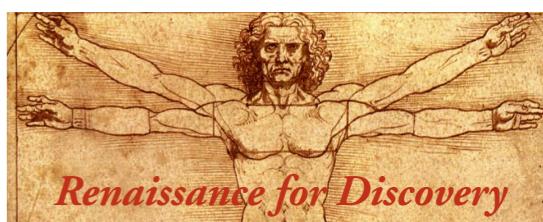
Excess at small q



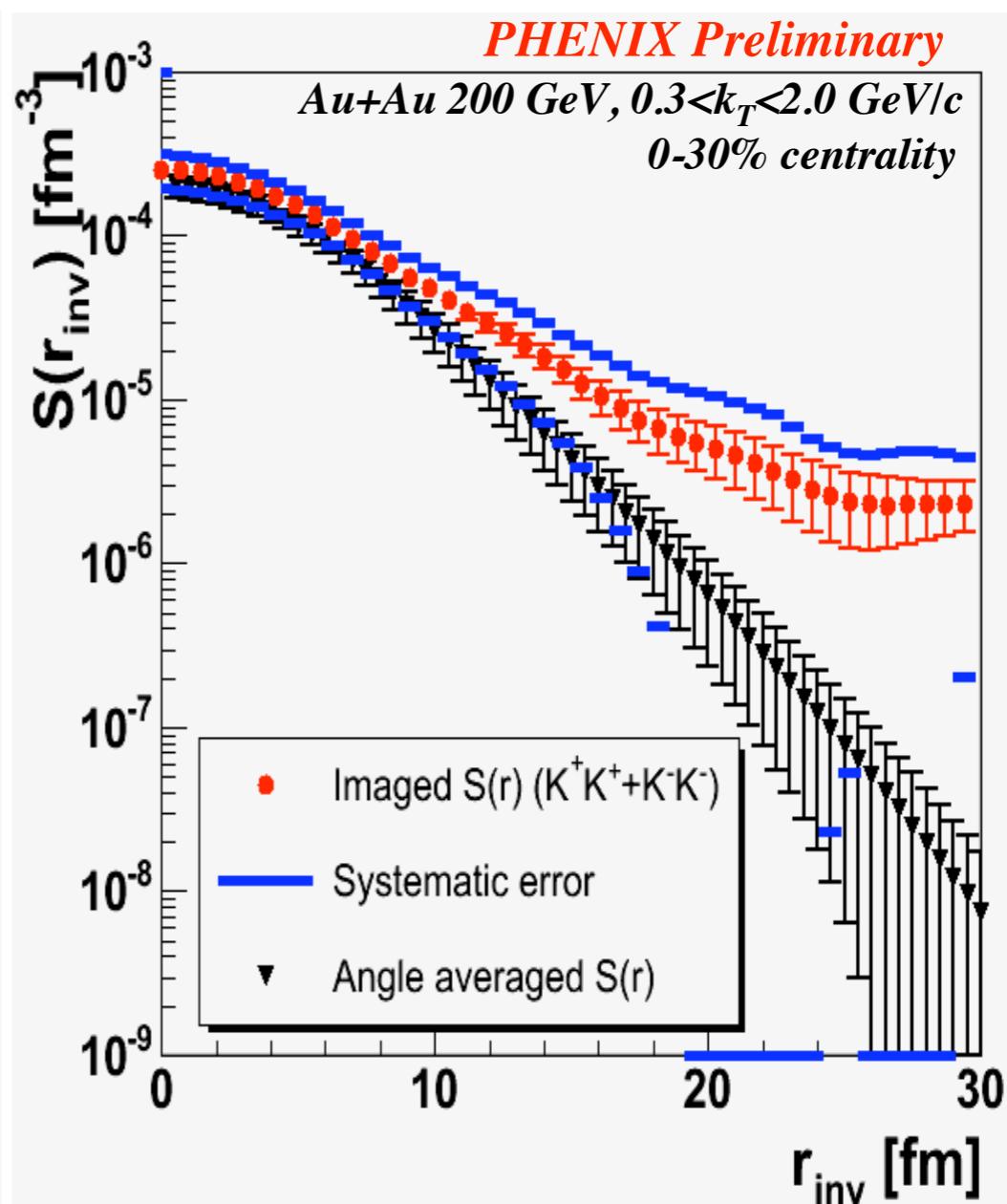
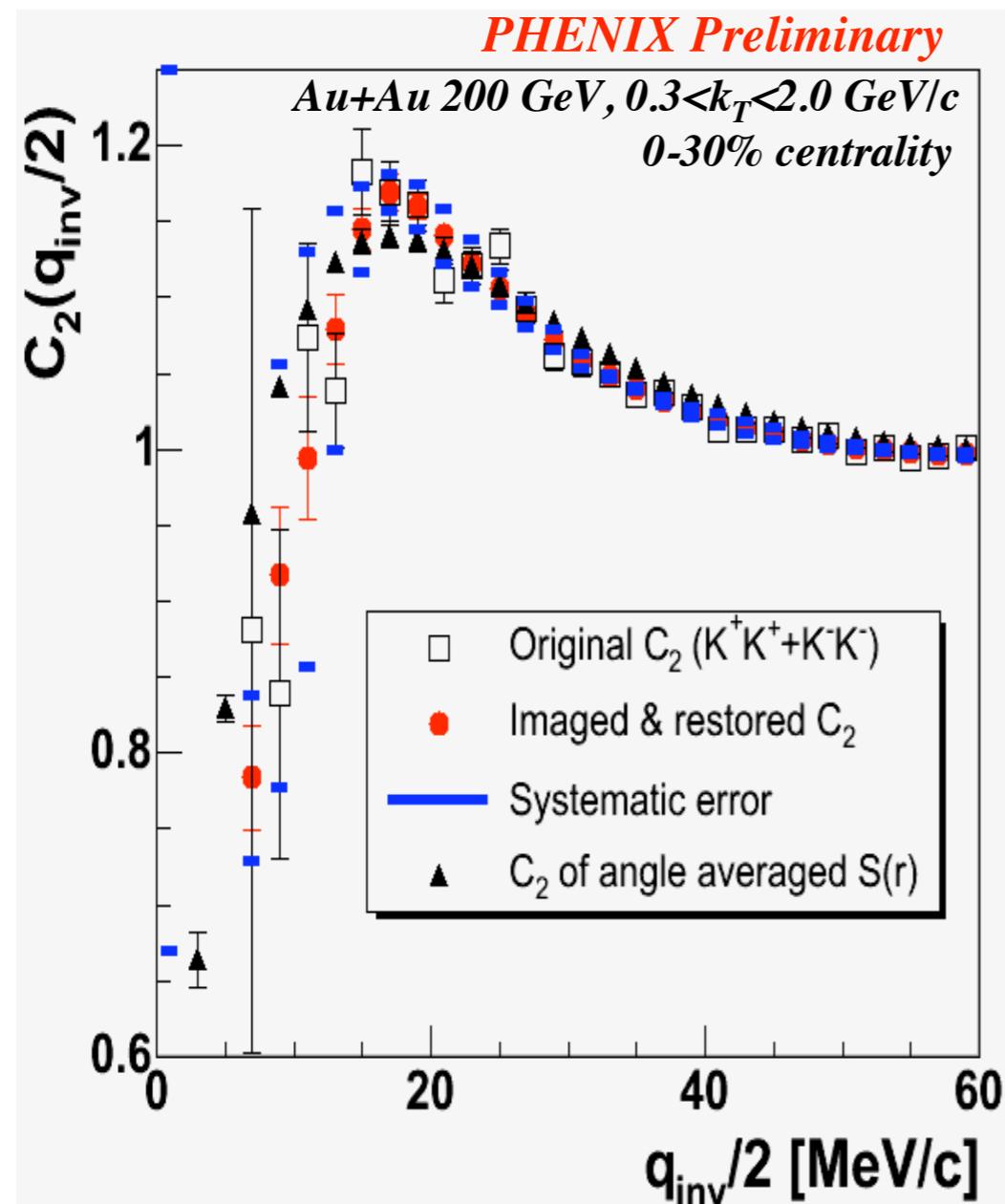
$$S(r) = \frac{\lambda R_{\text{eff}} \times e^{-\frac{r^2}{4R_T^2}} \operatorname{erfi}(\frac{r}{2R_{\text{eff}}})}{(8\pi R_T^2 R_0 r)}$$

Brown et al. PRC64(2001)014902

⇒ Excess at large R



Kaon ID Source Image



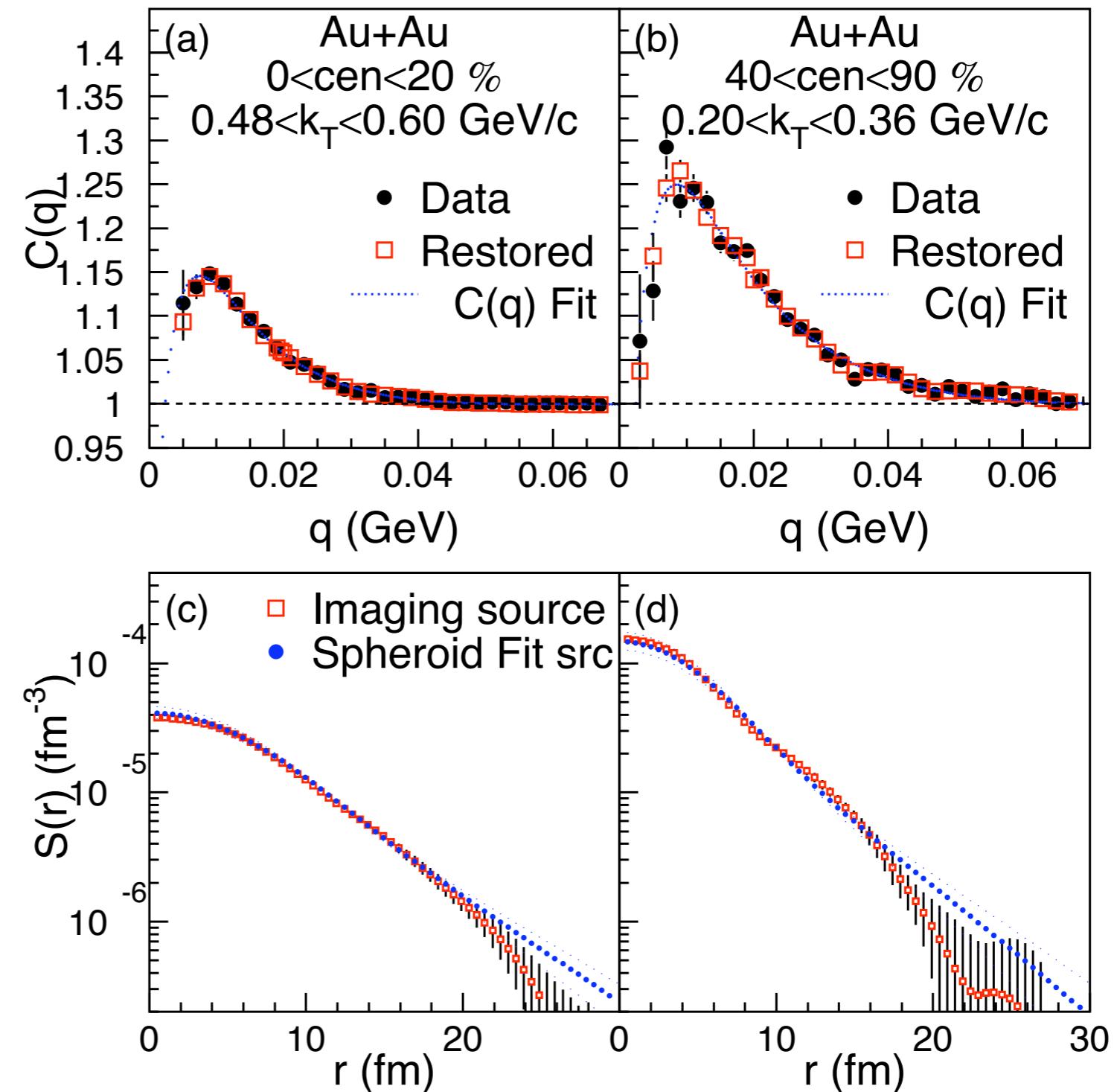
A. Enokizono, DNP Maui 2005

Kaon source image also has hint of long-range component.



Can we isolate the excess?

Long Range source not
equally present in all
kinematic regions.

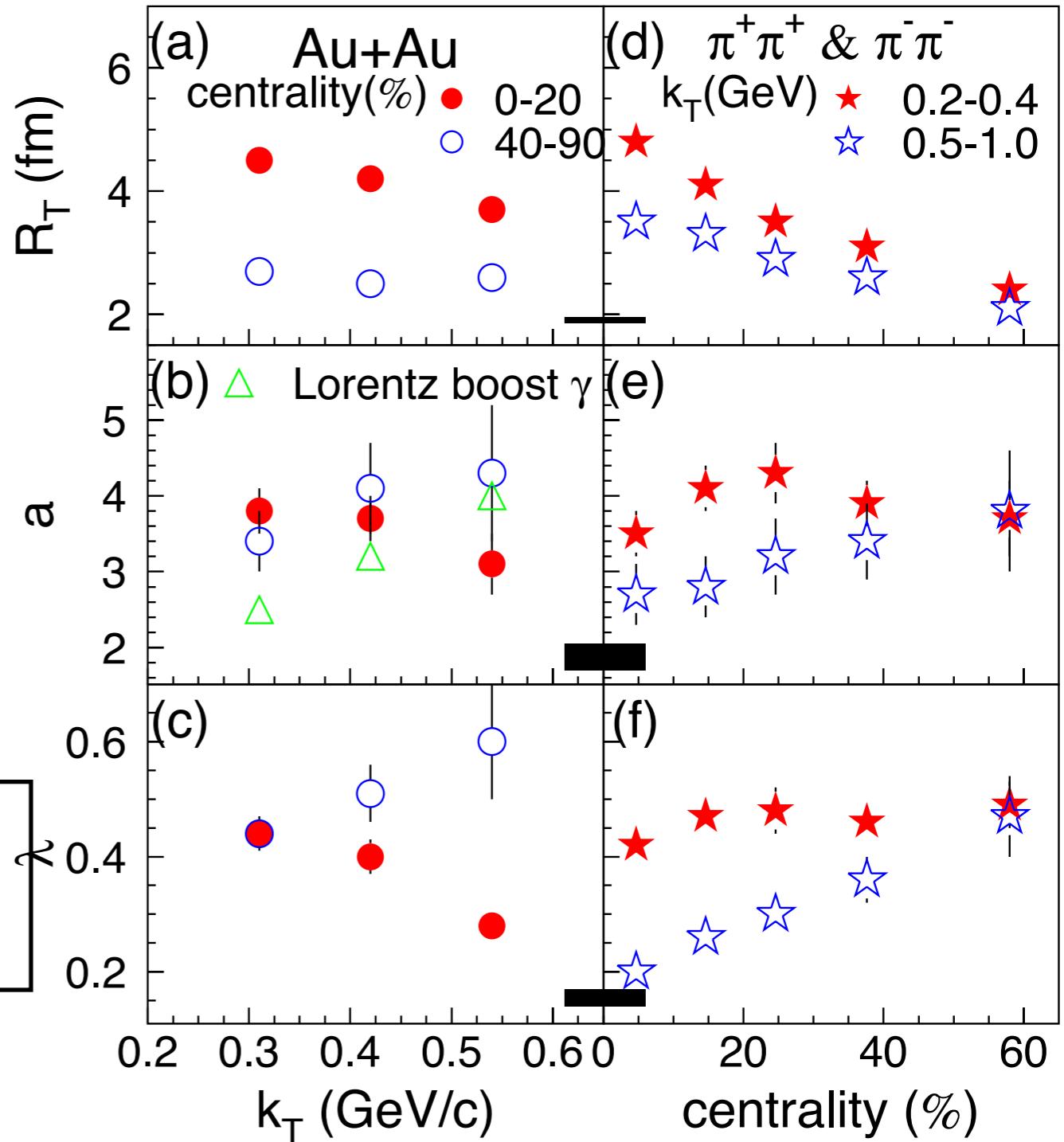


k_T , Centrality Systematics

Long Range contribution
most prominent at low- k_T
and in most central collisions.

$$R_O = a \times R_T$$

λ fraction of π pairs
contributing to the source



Summary

- PHENIX has measured two-particle correlations of π 's, K 's, and protons and three-particle correlations for π 's.
- 2π and $2K$ source radii demonstrate consistent $N_p^{1/3}$ scaling.
- 2π , $2K$, and $2p$ ID source Radii demonstrate consistent m_T scaling.
- 2π and $2K$ 3D source radii demonstrate consistent m_T scaling without indication of long emission duration nor significant hadronic rescattering.
- ID Imaging analyses of 2π indicate long range component that is most strongly present at low- k_T and central collisions.
- Preliminary ID Imaging analyses of $2K$ are also consistent with long range component.
- Images from Hydrodynamic calculations should be quantitatively compared with experimental images.

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August 1999

Department of Public Information
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*as of March 2005